Violation of the weak equivalence principle due to gravity-matter entanglement

Nikola Paunković^{1,2}, Francisco Pipa³ and Marko Vojinović⁴

- ¹ Department of Mathematics, IST, University of Lisbon
- 2 Security and Quantum Information Group (SQIG), Institute of Telecommunications, Lisbon
- ³ Department of Physics, IST, University of Lisbon
- ⁴ Group for Gravitation, Particles and Fields (GPF), Institute of Physics, University of Belgrade







Weak Equivalence Principle

The local effects of particle motion in a gravitational field are indistinguishable from those of an accelerated observer in flat spacetime.

Consequence:

A particle in a gravitational field should follow the geodesic, since this is how the straight line in flat space looks like from the accelerated frame.

Deriving WEP (geodesic motion) in GR

Single-pole approximation:

$$T^{\mu\nu}(x) = \int_{\mathcal{C}} d\tau \, B^{\mu\nu}(\tau) \frac{\delta^{(4)}(x - z(\tau))}{\sqrt{-g}} \,. \tag{1}$$

Conservation of stress-energy tensor (assuming the local Poincaré invariance for both $S_G[g]$ and $S_M[g,\phi]$):

$$\nabla_{\nu} T^{\mu\nu} = 0. \tag{2}$$

Replacing (2) into (1), we obtain the geodesic equation, with $u^{\mu} \equiv \frac{dz^{\mu}(\tau)}{d\tau}$ and $u^{\mu}u_{\mu} \equiv -1$ (Mathisson and Papapetrou [2, 3]; see also [4]):

$$u^{\lambda} \nabla_{\lambda} u^{\mu} = 0.$$

Using Cristoffel symbols,

$$\frac{d^2 z^{\lambda}(\tau)}{d\tau^2} + \Gamma^{\lambda}{}_{\mu\nu} \frac{dz^{\mu}(\tau)}{d\tau} \frac{dz^{\nu}(\tau)}{d\tau} = 0.$$

Quantising gravity

Fundamental gravitational degrees of freedom \hat{g} and $\hat{\pi}_g$:

$$\Delta \hat{g} \Delta \hat{\pi}_g \ge \frac{\hbar}{2}, \qquad \Delta \hat{\phi} \Delta \hat{\pi}_\phi \ge \frac{\hbar}{2}.$$

Separable state ($|g\rangle$ and $|\phi\rangle$ – coherent states of gravity and matter):

$$|\Psi\rangle = |g\rangle \otimes |\phi\rangle$$
.

Effective classical metric and stress-energy tensors:

$$g_{\mu\nu} \equiv \langle \Psi | \hat{g}_{\mu\nu} | \Psi \rangle , \qquad T_{\mu\nu} \equiv \langle \Psi | \hat{T}_{\mu\nu} | \Psi \rangle .$$

Violation of WEP due to entanglement

Entangled state (perturbation $|\tilde{\Psi}\rangle = |\tilde{g}\rangle \otimes |\tilde{\phi}\rangle$, with coherent classical states $|\tilde{g}\rangle$ and $|\tilde{\phi}\rangle$):

$$|\Psi\rangle = \alpha |\Psi\rangle + \beta |\tilde{\Psi}\rangle.$$

"Entangled" metric:

$$\boldsymbol{g}_{\mu\nu} = \langle \boldsymbol{\Psi} | \hat{g}_{\mu\nu} | \boldsymbol{\Psi} \rangle = g_{\mu\nu} + \beta h_{\mu\nu} + \mathcal{O}(\beta^2).$$

The perturbation is evaluated to be:

$$h_{\mu\nu} = 2 \operatorname{Re} \left[\langle \Psi | \hat{g}_{\mu\nu} | \tilde{\Psi} \rangle - \langle \Psi | \tilde{\Psi} \rangle g_{\mu\nu} \right] .$$

"Entangled" geodesic equation with the manifestly covariant correction:

$$\frac{d^2 z^{\mu}(\tau)}{d\tau^2} + \boldsymbol{\Gamma}^{\mu}{}_{\rho\nu} \frac{dz^{\rho}(\tau)}{d\tau} \frac{dz^{\nu}(\tau)}{d\tau} = 0,$$

$$u^{\lambda} \nabla_{\lambda} u^{\mu} + \beta \left(\nabla_{\rho} h^{\mu}{}_{\nu} - \frac{1}{2} \nabla^{\mu} h_{\nu\rho} \right) u^{\rho} u^{\nu} + \mathcal{O}(\beta^2) = 0.$$

Bibliography

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