VIOLATION OF THE WEAK EQUIVALENCE PRINCIPLE DUE TO GRAVITY-MATTER ENTANGLEMENT



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Abstract

We show the violation of the Weak Equivalence Principle in the presence of the entanglement between gravity and matter [1]. We analyse a simple toy scenario in which the entangled gravity-matter state consists of two terms: a dominant (classical) one, and a small perturbation. We rewrite the new geodesic equation as a geodesic equation for the dominant classical metric plus a contribution, coming from the interference terms between the dominant and perturbative gravity-matter fields. The additional term represents the deviation from the original geodesic equation due to the presence of gravity-matter entanglement, and measures the violation of the weak equivalence principle.

Weak Equivalence Principle (WEP)

The local effects of particle motion in a gravitational field are indistinguishable from those of an accelerated observer in flat spacetime.

Consequence

A particle in a gravitational field should follow the geodesic, since this is how the straight line in flat space looks like from the accelerated frame.

Single-pole approximation

Quantising gravity

Fundamental gravitational degrees of freedom \hat{g} and $\hat{\pi}_g$:

$$\Delta \hat{g} \Delta \hat{\pi}_g \ge \frac{\hbar}{2} \qquad \Delta \hat{\phi} \Delta \hat{\pi}$$

$$|\Psi\rangle = |g\rangle \otimes |\phi\rangle$$

 $|g\rangle$ and $|\phi\rangle$ – coherent states of gravity and matter.

 $g_{\mu\nu} \equiv \langle \Psi | \hat{g}_{\mu\nu} | \Psi \rangle \qquad T_{\mu\nu} \equiv \langle \Psi | \hat{T}_{\mu\nu} | \Psi$

$$T^{\mu\nu}(x) = \int_{\mathcal{C}} d\tau \, B^{\mu\nu}(\tau) \frac{\delta^{(4)}(x - z(\tau))}{\sqrt{-g}} \qquad (1)$$

Conservation of stress-energy tensor

We assume the local Poincaré invariance for both $S_G[g]$ and $S_M[g, \phi]$

 $\nabla_{\nu}T^{\mu\nu} = 0 \tag{2}$

Derivation of WEP (geodesic motion) from General Relativity

Replacing (2) into (1), we obtain the geodesic equation, with $u^{\mu} \equiv \frac{dz^{\mu}(\tau)}{d\tau}$ and $u^{\mu}u_{\mu} \equiv -1$ (Mathisson and Papapetrou [2, 3]; see also [4]) $u^{\lambda} \nabla_{\lambda} u^{\mu} = 0$
$$\begin{split} \textbf{Violation of WEP due to entanglement} \\ \hline \textbf{Fintangled" metric} \\ \textbf{g}_{\mu\nu} = \langle \Psi | \hat{g}_{\mu\nu} | \Psi \rangle \\ \hline \textbf{W} \rangle = \alpha | \Psi \rangle + \beta | \tilde{\Psi} \rangle \\ \hline \textbf{Ferturbation, with coherent} \\ \textbf{chassical states } | \tilde{g} \rangle & \text{and } | \tilde{g} \rangle \\ \hline \textbf{W} \rangle = g \mu \nu + \beta h \mu \nu + \beta (\beta^2) \\ \hline \textbf{W} \rho = 2 \operatorname{Re} \left[\langle \Psi | \hat{g}_{\mu\nu} | \tilde{\Psi} \rangle - \langle \Psi | \tilde{\Psi} \rangle g_{\mu\nu} \right] \\ \mu \nu = 2 \operatorname{Re} \left[\langle \Psi | \hat{g}_{\mu\nu} | \tilde{\Psi} \rangle - \langle \Psi | \tilde{\Psi} \rangle g_{\mu\nu} \right] \\ \mu^{\lambda} \nabla_{\lambda} u^{\mu} + \beta \left(\nabla_{\rho} h^{\mu}{}_{\nu} - \frac{1}{2} \nabla^{\mu} h_{\nu\rho} \right) u^{\rho} u^{\nu} + \mathcal{O}(\beta^2) = 0 \end{split}$$

$$\frac{d^2 z^{\mu}(\tau)}{d\tau^2} + \boldsymbol{\Gamma}^{\mu}{}_{\rho\nu} \frac{dz^{\rho}(\tau)}{d\tau} \frac{dz^{\nu}(\tau)}{d\tau} = 0 , \qquad \qquad$$

Discussion

- Both matter and gravity are considered quantum no semiclassical approximations and the associated errors.
- In Newtonian physics, due to the specific dynamical and gravitational laws $(m_i a = m_g g)$, WEP implies the equality of two types of masses, $m_i = m_g$. The laws of quantum mechanics, however, may not directly imply that $m_i = m_g$ as a consequence of WEP. On the other hand, the deviation from the geodesic trajectory is a more universal signal of WEP violation.
- Further analysis of (2) for the 'entangled" stress-energy tensor $T_{\mu\nu} = \langle \Psi | \hat{T}_{\mu\nu} | \Psi \rangle$ can be done obtaining the domain of validity of the single-pole approximation.
- Quantitative analysis might allow for possible future experimental verification of the gravity-matter entanglement and thus an answer to the open question of the necessity of quantising gravity.

where the Cristoffel symbols are given by:

 $\boldsymbol{\Gamma}^{\mu}{}_{\rho\nu} = \frac{1}{2} \boldsymbol{g}^{\mu\sigma} (\partial_{\rho} \boldsymbol{g}_{\nu\sigma} + \partial_{\nu} \boldsymbol{g}_{\sigma\rho} - \partial_{\sigma} \boldsymbol{g}_{\rho\nu}) \,.$

References

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