# SPINCUBE MODEL OF QG AND CONNECTION TO CDT

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### CONCEPTS

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Nature has a physical cutoff at the scale  $\gamma l_p$ .

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The simplest implementation of the principle:

# Spacetime is a piecewise-linear manifold.

Things to note:

- PL-structure (a triangulation) is not a regulator, but a physical entity;
- UV-completion: inside a 4-simplex, spacetime is flat and matter fields are constant;
- finite number of degrees of freedom (in a finite volume);
- field theory reconstructed only as an approximation, like in fluid mechanics.

Classical theory — constrained BFCG action:

$$S = \int \underbrace{B^{ab} \wedge R_{ab} + e^a \wedge \nabla \beta_a}_{\text{topological theory}} - \underbrace{\phi^{ab} \wedge \left(B_{ab} - \varepsilon_{abcd} \ e^c \wedge e^d\right)}_{\text{simplicity constraint}}.$$

Basic properties:

- equivalent to general relativity,
- similar in structure to the Plebanski action,
- contains tetrad fields in the topological sector,
- coupling to matter fields completely straightforward,
- based on 2BF action for the Poincaré 2-group.

object	symbol	spinfoams	spincube
vertex	v		
edge	$\epsilon$		$l \in \mathbb{R}_0^+$
triangle	$\Delta$	$j \in \mathbb{N}_0/2$	$m \in \mathbb{Z}$
tetrahedron	au	$\iota \in \mathbb{N}_0/2$	M = 1
4-simplex	$\sigma$		

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Quantum theory — the spincube state sum:

$$Z = \sum_{T \in \mathcal{T}} \int_{\mathbb{R}^+_0} dl_1 \dots \int_{\mathbb{R}^+_0} dl_E \sum_{m_1 \in \mathbb{Z}} \dots \sum_{m_F \in \mathbb{Z}} \prod_{\epsilon \in T} \mathcal{A}_{\epsilon}(l,m) \prod_{\Delta \in T} \mathcal{A}_{\Delta}(l,m) \prod_{\sigma \in T} e^{iS_{\sigma}(l,m)} \mathcal{A$$

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Edge and triangle amplitudes  $\mathcal{A}_{\epsilon}$  and  $\mathcal{A}_{\Delta}$  are chosen such that they impose the simplicity constraint between *l*'s and *m*'s:

$$|m_{\Delta}| = \frac{1}{\gamma l_p^2} A_H(l_{\epsilon_1}, l_{\epsilon_2}, l_{\epsilon_3}), \qquad \epsilon_1, \epsilon_2, \epsilon_3 \in \Delta, \qquad \forall \Delta \in T.$$

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$$|m_{1}| = A_{1}(l_{1}, \dots, l_{E}), \\ \vdots \\ |m_{E}| = A_{E}(l_{1}, \dots, l_{E}), \\ |m_{E+1}| = A_{E+1}(l_{1}, \dots, l_{E}), \\ \vdots \\ |m_{F}| = A_{F}(l_{1}, \dots, l_{E}), \end{cases} \Rightarrow \begin{cases} l_{1} = L_{1}(m_{1}, \dots, m_{E}), \\ l_{E} = L_{E}(m_{1}, \dots, m_{E}), \\ l_{E} = L_{E}(m_{1}, \dots, m_{E}), \\ |m_{E+1}| = f_{1}(m_{1}, \dots, m_{E}), \\ \vdots \\ |m_{F}| = f_{F-E}(m_{1}, \dots, m_{E}). \end{cases}$$

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$$\begin{aligned} |m_1| &= A_1(l_1, \dots, l_E), \\ \vdots &= \\ |m_E| &= A_E(l_1, \dots, l_E), \\ |m_{E+1}| &= A_{E+1}(l_1, \dots, l_E), \\ \vdots &= \\ |m_F| &= A_F(l_1, \dots, l_E), \end{aligned} \right\} \Rightarrow \begin{cases} l_1 &= L_1(m_1, \dots, m_E), \\ l_E &= L_E(m_1, \dots, m_E), \\ |m_{E+1}| &= f_1(m_1, \dots, m_E), \\ \vdots &= \\ |m_F| &= f_{F-E}(m_1, \dots, m_E) \end{aligned}$$

Either impose the simplicity constraint weakly (i.e. on-shell, see talk by A. Miković), or prove that the system of Diophantine equations has a nonempty set of solutions!

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$$\begin{cases} 4 \\ 5 \\ 1 \\ 1 \\ 1 \\ 1 \\ 2 \\ \end{cases}^{3} \qquad |m_{123}| = A_{H}(l_{12}, l_{13}, l_{23}) \\ \vdots \\ |m_{345}| = A_{H}(l_{34}, l_{35}, l_{45}) \\ \end{cases} \Rightarrow \begin{cases} l_{12} = L_{12}(m_{123}, \dots, m_{345}), \\ \vdots \\ l_{45} = L_{45}(m_{123}, \dots, m_{345}). \end{cases}$$

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The case of general triangulation  $(F \ge E)$ : proof by induction! In the generic case, all 4-simplices in the triangulation must be equal.

In general, given a solution of type

$$l_{\epsilon} = L_{\epsilon}(m_1, \dots, m_{10}), \qquad \forall \epsilon \in T,$$

we implement the simplicity constraint by choosing the amplitudes as

$$\mathcal{A}_{\epsilon}(l,m) = \delta(l_{\epsilon} - L_{\epsilon}), \qquad \mathcal{A}_{\Delta}(l,m) = \chi\left(|m_{\Delta}| - \frac{1}{\gamma l_{p}^{2}}A_{H}(L_{\epsilon \in \Delta})\right),$$

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so the spincube state sum becomes:

$$Z = \sum_{T \in \mathcal{T}} \sum_{m_1 \in \mathbb{Z}} \dots \sum_{m_{10} \in \mathbb{Z}} e^{iS_R(L(m))}$$

Since  $m_1, \ldots, m_{10}$  belong to the same 4-simplex, we can commute the sums and write

$$Z = \sum_{m_1 \in \mathbb{Z}} \dots \sum_{m_{10} \in \mathbb{Z}} \left( \sum_{T \in \mathcal{T}} e^{iS_R(L(m))} \right).$$

Consider the special case of isosceles 4-simplices, such that  $l_{\epsilon} \in \{a, b\}$ . There are 40 such simplices in total:



plus the "thick $\leftrightarrow$ thin" inverted ones ("thin" = a, "thick" = b). In general, these simplices contain four types of triangles,

 $(a,a,a), \qquad (a,a,b), \qquad (a,b,b), \qquad (b,b,b),$ 

so the simplicity constraint consists of 4 equations for 2 variables, and has no solution.

But three pairs of simplices contain only two types of triangles,



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For these the simplicity constraint reduces to the system of two equations for two variables, and can always be solved. The state sum can then be written as:

$$Z = \sum_{m_1 \in \mathbb{Z}} \sum_{m_2 \in \mathbb{Z}} \left( \sum_{T \in \mathcal{T}} e^{iS_R(a(m), b(m))} \right)$$

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$$Z = \sum_{m_1 \in \mathbb{Z}} \sum_{m_2 \in \mathbb{Z}} \underbrace{\left( \sum_{T \in \mathcal{T}} e^{iS_R(a(m), b(m))} \right)}_{Z_{CDT}}.$$

# CONCLUSIONS

### Overview of results:

- simplicity constraint can be imposed strongly, so that the spincube model is well defined, despite overcomplete system of equations;
- two classes of solutions: identical irregular 4-simplices, and three types of isosceles 4-simplices;
- CDT is one of isosceles solutions, and thus a special case of the spincube model.

#### Topics for further research:

- classify solutions containing only  $3, 4, \ldots, 9$  different edges;
- repeat the CDT phase-transition analysis for all other classes, study what happens to phases when one varies the edge-length parameters;
- introduce matter fields, study realistic QG systems;
- study the relationship between the two semiclassical limits: average over many triangulations, or take  $l, m \to \infty$ , or both?
- $\Rightarrow$  More people is needed to do all this!

### THANK YOU!